

SOLVABLE VARIATIONAL PROBLEMS IN NON EQUILIBRIUM STATISTICAL MECHANICS

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The general Framework

Non equilibrium models of statistical mechanics

Presence of fluxes (matter, energy,...) in the stationary state

Non reversible stochastic models

Large deviations analysis \implies Variational problems

Large deviations

X_N sequence of random variables taking values on a metric space M satisfies a Large deviations Principle (**LDP**) with speed $\alpha(N)$ and rate $I : M \rightarrow \mathbb{R}^+ \cup \{+\infty\}$ if

$$\limsup_{N \rightarrow +\infty} \frac{1}{\alpha(N)} \log \mathbb{P}(X_N \in C) \leq - \inf_{x \in C} I(x), \quad \forall C \text{ closed}$$

$$\liminf_{N \rightarrow +\infty} \frac{1}{\alpha(N)} \log \mathbb{P}(X_N \in O) \geq - \inf_{x \in O} I(x), \quad \forall O \text{ open}$$

$$\mathbb{P}(X_N \sim x) \simeq e^{-\alpha(N)I(x)}$$

Examples

- $X_N \in \mathbb{R}^d$ Gaussian

$$\frac{1}{Z_N} e^{-\frac{N}{2}(x, C^{-1}x)} dx$$

satisfies a LDP with speed N and rate

$$I(x) = \frac{1}{2} (x, C^{-1}x)$$

- A gas of independent coins

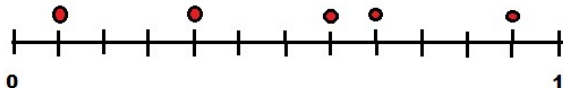
$$\eta = \left(\eta(1), \dots, \eta(N) \right) \in \{0, 1\}^N, \quad \text{i.i.d. Bernoulli}$$

Examples

EMPIRICAL MEASURE

$$\eta \implies \pi_N(\eta) := \frac{1}{N} \sum_{i=1}^N \eta(i) \delta_{\frac{i}{N}}$$

lattice size = $1/N$



● = delta measure

LD for the empirical measure

$\pi_N(\eta) \in \mathcal{M}^+([0, 1])$, random positive measure, satisfies LDP

$$\mathbb{P}(\pi_N(\eta) \sim \rho(x)dx) \simeq e^{-NI(\rho)}$$

the rate function is

$$I(\rho) = \int_0^1 h_m(\rho(x)) dx$$

where

$$h_m(\rho) = \rho \log \frac{\rho}{m} + (1 - \rho) \log \frac{1 - \rho}{1 - m}$$

Contraction

Average number of coins

$$\frac{1}{N} \sum_i \eta(i) = \int_{[0,1]} d\pi_N(\eta)$$

satisfies LDP

$$\mathbb{P} \left(\frac{1}{N} \sum_i \eta(i) \sim y \right) \simeq e^{-NJ(y)}$$

BY CONTRACTION

$$J(y) = \inf_{\{\rho: \int_0^1 \rho(x) dx = y\}} I(\rho)$$

A simple non equilibrium model

CONTINUOUS TIME MARKOV CHAINS WITH "EXPONENTIALLY SMALL" RATES OF TRANSITION

V finite set $x, y, z \in V$; $N \sim \frac{1}{T}$, $T = \text{temperature}$

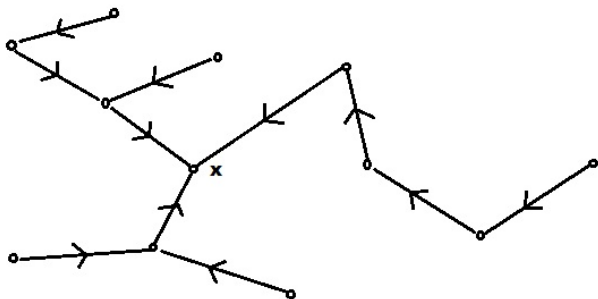
TRANSITION RATES

$$r_N(y, z) = e^{-NR(y,z)}$$

INVARIANT MEASURE

$$\mu_N(x) = \frac{\sum_{g \in G_x} \prod_{(y,z) \in g} r_N(y,z)}{Z_N}$$

The graphs in \mathbb{G}_x



A simple non equilibrium model

$$\mu_N(x) \simeq e^{-NW(x)}$$

$$W(x) = \min_{\mathcal{G} \in \mathcal{G}_x} \left\{ \sum_{(y,z) \in \mathcal{G}} R(y,z) \right\} + \text{constant}$$

The reversible case (Example)

$$R(y,z) = \frac{H(z) - H(y)}{2}$$

Combinatorial optimization problem is solved (tricky)

$$W = H + \text{constant}$$

From dynamic to static LD

Stochastic differential equation in \mathbb{R}^d

$$dX^\epsilon(t) = b(X^\epsilon(t)) dt + \sqrt{\epsilon} dW(t)$$

Dynamic sample path LD

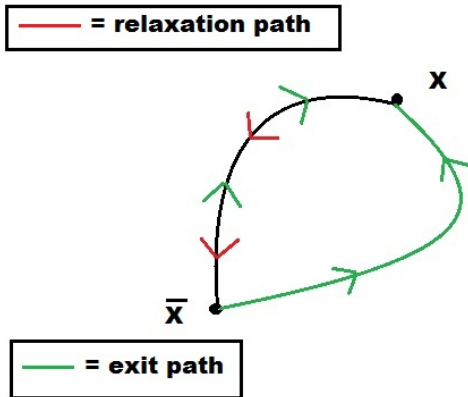
$$\mathbb{P}\left(X^\epsilon(t) \sim x(t), t \in [T_1, T_2]\right) \simeq e^{-\frac{\epsilon^{-1}}{2} \int_{T_1}^{T_2} |\dot{x} - b(x)|^2 dt}$$

$\left\{ \begin{array}{l} b = \text{globally attractive vector field} \\ \bar{x} \text{ such that } b(\bar{x}) = 0 \text{ unique equilibrium point} \end{array} \right.$

Invariant measure μ^ϵ satisfies LD with rate the QUASIPOTENTIAL

$$V(x) = \inf_{T>0} \inf_{\{x(t) : x(-T)=\bar{x}, x(0)=x\}} I_{[-T,0]}(x(t))$$

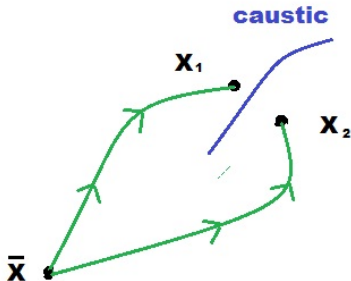
The quasipotential



The quasipotential

$$b = -\nabla S \iff \text{reversible} \implies V = 2S$$

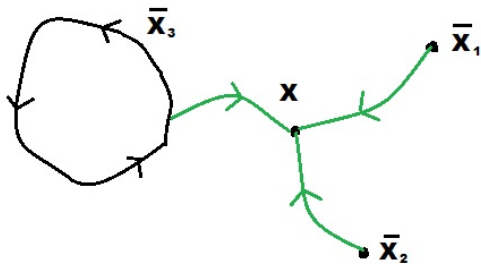
In general V not differentiable (phase transitions, WASEP)



The general case

If b has several stable attractors the quasipotential becomes

$$V(x) = \inf_i W_i + V_i(x)$$



A solvable case

A. Faggionato, D.G. (2012)

Stochastic differential equation on \mathbb{S}^1 (unit circle)

$$dX^\epsilon(t) = b(X^\epsilon(t)) dt + \sqrt{\epsilon} dW(t)$$

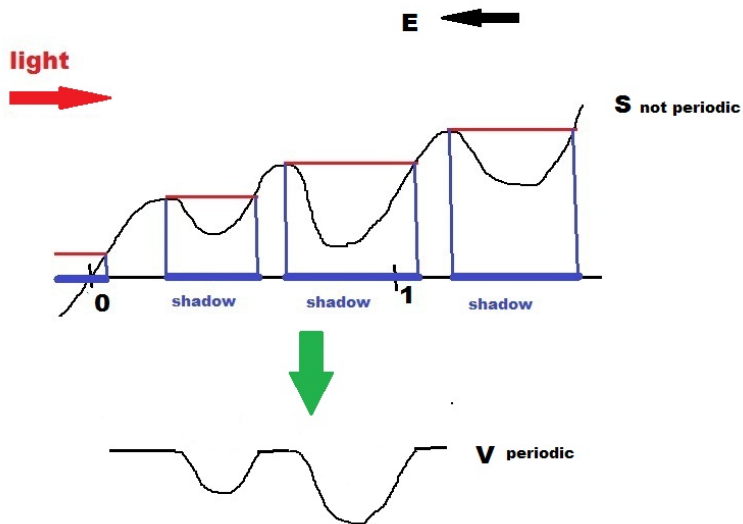
b is periodic of period one

$$S(x) := -2 \int_0^x b(y) dy$$

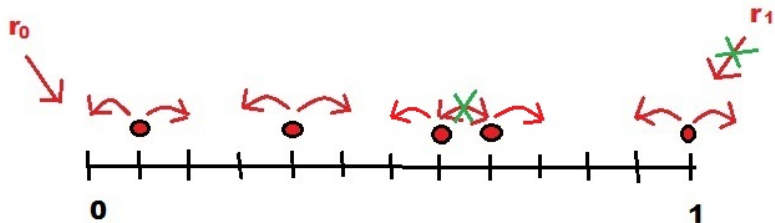
$$S \text{ periodic} \iff \int_0^1 b(y) dy = 0 \iff \text{reversible} \implies V = S$$

Add an external field $b \implies b + E$

Sunshine transformation



Boundary driven 1-d simple exclusion



Scaling limit

Diffusive rescaling

$$L_N \implies N^2 L_N$$

Hydrodynamic scaling limit

$$\pi_N(\eta_t) \xrightarrow{N \rightarrow +\infty} \rho(x, t) dx$$

$$\begin{cases} \rho_t = \Delta \rho \\ \rho(0, t) = \rho_{r_0} \\ \rho(1, t) = \rho_{r_1} \end{cases}$$

Large deviations and quasipotential

BDGJL (2002)

DYNAMIC LARGE DEVIATIONS

$$\mathbb{P}(\pi_N(\eta_t) \sim \rho(x, t) dx, t \in [T_1, T_2]) \simeq e^{-NI_{[T_1, T_2]}(\rho)}$$

$$\begin{cases} I_{[T_1, T_2]}(\rho) = \frac{1}{4} \int_{T_1}^{T_2} dt \int_0^1 dx \rho(1 - \rho) (\nabla H)^2 \\ \partial_t \rho = \Delta \rho - \nabla \cdot (\rho(1 - \rho) \nabla H), \quad H(0, t) = H(1, t) = 0 \end{cases}$$

THE QUASIPOTENTIAL

$r_0 = r_1 \implies$ reversible , gas of independent coins

$r_0 \neq r_1 \implies$ not reversible, long range correlations

$$\begin{cases} V(\rho) = \int_0^1 [h_f(\rho) + \log \frac{\nabla f}{\nabla \rho}] dx \\ \frac{\Delta f}{(\nabla f)^2} f(1 - f) + f = \rho, \quad f(0) = \rho_{r_0}, f(1) = \rho_{r_1} \end{cases}$$

The minimization path

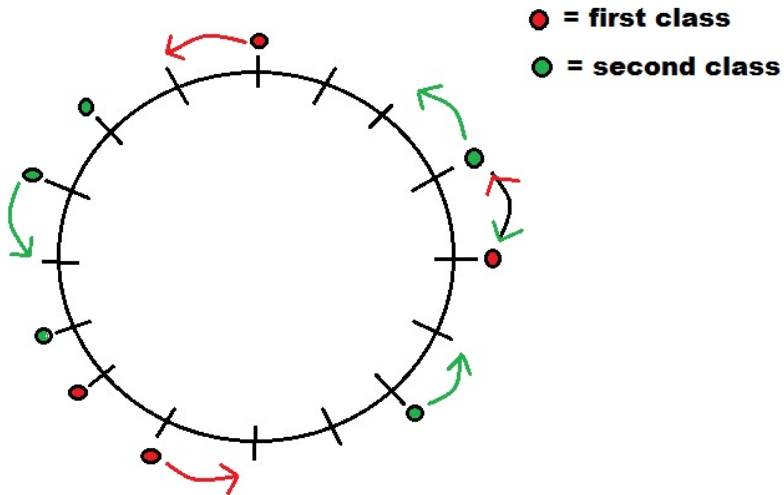
The minimizer for $V(\rho_0)$ is the time reversal of the following coupled differential problem

$$\begin{cases} \partial_t \rho = \Delta \rho - \nabla \cdot \left(\frac{\rho(1-\rho)}{f(1-f)} \nabla f \right) \\ \rho(u, 0) = \rho_0(u) \\ f(1-f) \frac{\Delta f}{(\nabla f)^2} + f = \rho \\ \text{and b.c.} \end{cases}$$

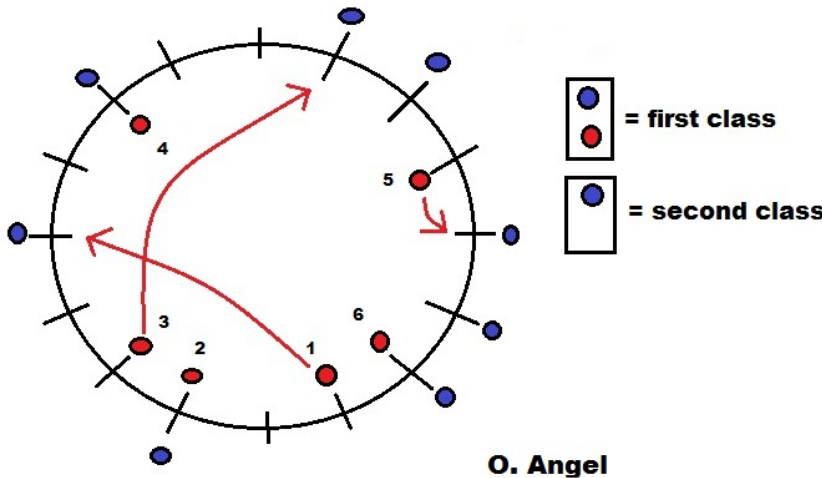
A computation (magic transformation!) shows it is equivalent to

$$\begin{cases} \partial_t f = \Delta f \\ f(1-f) \frac{\Delta f}{(\nabla f)^2} + f = \rho \\ \rho(u, 0) = \rho_0(u) \\ \text{and b.c.} \end{cases}$$

2-class TASEP



The invariant measure



Collapsing particles

$$(\tilde{\eta}_1, \tilde{\eta}_T) : \sum_x \tilde{\eta}_1(x) \leq \sum_x \tilde{\eta}_T(x) \implies (\eta_1, \eta_T) = \mathcal{C}[(\tilde{\eta}_1, \tilde{\eta}_T)]$$

Flux across bond $(x, x + 1)$

$$J(x) = \sup_y \left[\sum_{z \in [y, x]} \tilde{\eta}_1(z) - \tilde{\eta}_T(z) \right]_+$$

Collapsing measures

D.G. (08)

$$(\tilde{\rho}_1, \tilde{\rho}_T) : \int_{\mathbb{S}^1} d\tilde{\rho}_1 \leq \int_{\mathbb{S}^1} d\tilde{\rho}_T \implies (\rho_1, \rho_T) = \mathcal{C}[(\tilde{\rho}_1, \tilde{\rho}_T)]$$

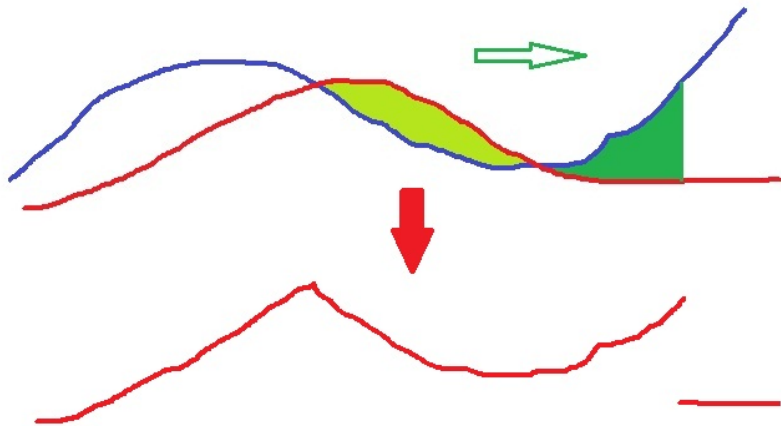
Definition

$$\int_{(a,b]} d\rho_1 = \int_{(a,b]} d\tilde{\rho}_1 + J(a) - J(b)$$

where

$$J(x) := \sup_y \left[\int_{(y,x]} d\tilde{\rho}_1 - \int_{(y,x]} d\tilde{\rho}_2 \right]_+$$

Collapsing measures



Large deviations

LD for the $(\tilde{\eta}_1, \tilde{\eta}_T)$ variables

$$\tilde{V}(\tilde{\rho}_1, \tilde{\rho}_T) = \int_{\mathbb{S}^1} [h_{m_1}(\tilde{\rho}_1) + h_{m_2}(\tilde{\rho}_T)] dx$$

LD for the SNS (not convex!)

$$\begin{aligned} V(\rho_1, \rho_T) &= \inf_{\{(\tilde{\rho}_1, \tilde{\rho}_T) : \mathcal{C}[(\tilde{\rho}_1, \tilde{\rho}_T)] = (\rho_1, \rho_T)\}} \tilde{V}(\tilde{\rho}_1, \tilde{\rho}_T) \\ &= \int_{\mathbb{S}^1} [h_{m_1}(\hat{\rho}_1) + h_{m_2}(\rho_T)] dx \end{aligned}$$

On any (a, b) where $\rho_1 = \rho_T$

$$\int_a^x \hat{\rho}_1(y) dy = CE \left[\int_a^x \rho_1(y) dy \right]$$