

Contact mechanics in an affine nutshell

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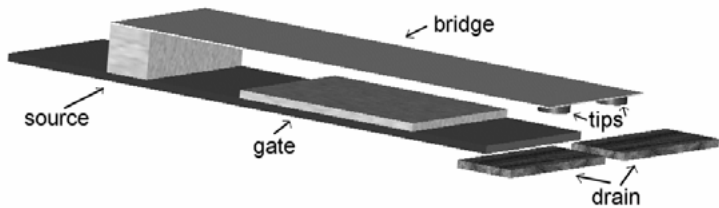
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Abstract

A nice tutorial in contact mechanics is readily obtained by considering an incompressible Mooney-Rivlin affine body undergoing short-range interactions with a flat, rigid surface – interactions meant to model impact, damping, friction and adhesion. Numerical integration of a handful of ODE's provides vivid simulations of the motion of an affine body – stiff or soft, heavy or light – that bounces, rocks, rolls, slides or gets stuck on a flat floor or a vertical wall.



Outline

- ▶ Rigid body
- ▶ Affine body
- ▶ Equations of motion
- ▶ Piola stress
 - ▶ Mooney-Rivlin material
 - ▶ Incompressibility
 - ▶ Dissipation
- ▶ Contact forces constitutive laws
 - ▶ Repulsion
 - ▶ Damping
 - ▶ Friction
 - ▶ Adhesion
- ▶ Bouncing, sliding, rolling, etc. (numerical simulations)

Rigid body

A motion of the body \mathcal{B} is described at each time t by a placement defined on the *paragon shape* \mathcal{D} :

$$p : \mathcal{D} \times \mathcal{J} \rightarrow \mathcal{E}$$

characterized by the following representation:

$$p(x, t) = p_o(t) + R(t)(x - x_o)$$

where $R(t) : \mathcal{V} \rightarrow \mathcal{V}$ is a rotation in the translation space of \mathcal{E} .

Test velocity fields

$$w(x) = w_o + W R(t)(x - x_o), \quad \text{with} \quad \text{sym } W = 0$$

Rigid body

Balance principle

$$\int_{\mathcal{D}} b \cdot w \, dV + \int_{\partial\mathcal{D}} q \cdot w \, dA = 0 \quad \forall w, \forall t$$

Equations of motion

$$-m \ddot{p}_o(t) - m g + f(t) = 0$$

$$\text{skw}(-\ddot{R}(t) J R(t)^T + M(t) R(t)^T) = 0$$

$$m := \int_{\mathcal{D}} \rho \, dV; \quad J := \int_{\mathcal{D}} \rho (x - x_o) \otimes (x - x_o) \, dV$$

$$f(t) := \sum \int_{\partial\mathcal{D}} q_j(x, t) \, dA; \quad M(t) := \sum \int_{\partial\mathcal{D}} (x - x_o) \otimes q_j(x, t) \, dA$$

Affine body

A motion of the body \mathcal{B} is described at each time t by a placement defined on the *paragon shape* \mathcal{D} :

$$p : \mathcal{D} \times \mathcal{J} \rightarrow \mathcal{E}$$

characterized by the following representation:

$$p(x, t) = p_o(t) + F(t)(x - x_o)$$

where $F(t) : \mathcal{V} \rightarrow \mathcal{V}$ is linear and such that $\det F(t) > 0$

Test velocity fields

$$w(x) = w_o + G F(t)(x - x_o)$$

Affine body

Balance principle

$$-S \cdot GF \operatorname{vol} \mathcal{D} + \int_{\mathcal{D}} b \cdot w \, dV + \int_{\partial \mathcal{D}} q \cdot w \, dA = 0 \quad \forall w, \forall t$$

Equations of motion

$$-m \ddot{p}_o(t) - mg + f(t) = 0$$

$$-\ddot{F}(t) J F(t)^T + (M(t) - S(t) \operatorname{vol} \mathcal{D}) F(t)^T = 0$$

Piola stress and material properties

Frame indifference

$$S \cdot WF = 0 \quad \forall W \mid \text{sym } W = 0 \quad \Rightarrow \quad \text{skw } SF^T = 0$$

Mooney-Rivlin strain energy (incompressible material)

$$\varphi(F) := c_{10}(\iota_1(C) - 3) + c_{01}(\iota_2(C) - 3)$$

Stress response (energetic + reactive + dissipative)

$$SF^T = \hat{S}(F)F^T - \pi I + \mu \dot{F}F^{-1}$$

$$\hat{S}(F) \cdot \dot{F} = d\varphi(F)/dt$$

$$\Rightarrow \quad \hat{S}(F)F^T = 2(c_{10}FF^T - c_{01}F^{-T}F^{-1})$$

Dissipation principle

$$S \cdot \dot{F} - d\varphi(F)/dt \geq 0 \quad \Rightarrow \quad \mu \geq 0$$

Contact force constitutive laws

Repulsive force

$$q_r(x, t) = \alpha_r d(x, t)^{-\nu_r} e_2$$

Damping force

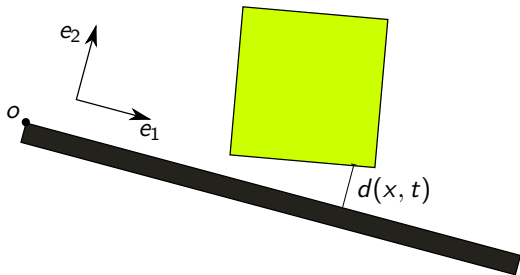
$$q_d(x, t) = -\alpha_d d(x, t)^{-\nu_d} (\dot{p}(x, t) \cdot e_2) e_2$$

Frictional force

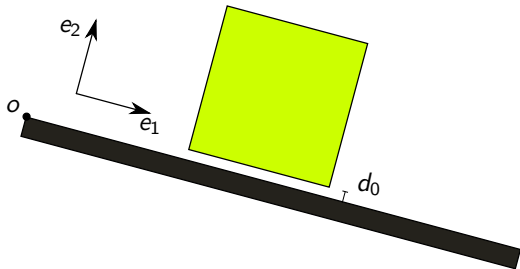
$$q_f(x, t) = -\alpha_f d(x, t)^{-\nu_f} (\dot{p}(x, t) \cdot e_1) e_1$$

Adhesive force

$$q_a(x, t) = -\alpha_a d(x, t)^{-\nu_a} e_2$$



$$d(x, t) := ((p(x, t) - o) \cdot e_2)$$



$$d(x, t) := ((p(x, t) - o) \cdot e_2) e_2$$

Contact force working

The contact forces are surface forces per unit deformed area:

$$\int_{\partial\mathcal{D}} q(x, t) \cdot w(x, t) k(x, t) dA$$

Area change factor:

$$k(x, t) := \|F(t)^{-T} n(x)\| \det(F(t))$$

$n(x)$ outward unit normal vector on $\partial\mathcal{D}$

Numerical simulations

rigid block; contact: repulsion, damping, no friction

001 | 002 | 003

rigid block; contact: repulsion, damping, friction

011 | 012 | 013

rigid block on a sloping plane

021 | 022 | 023

rigid disk

031 | 032 | 033 | 034 | 035

affine disk

041 | 042

Appendix

Cauchy stress

$$T = SF^T \frac{1}{\det F}$$

Pressure π

It is the *reactive* part of T . In an incompressible solid/fluid the velocity fields are said to be *isochoric*. The trace of the velocity gradient turns out to be zero.

A reactive stress, whose power is zero for any isochoric velocity field, has to be a spherical tensor $-\pi I$:

$$\pi I \cdot G = \pi \operatorname{tr} G = 0$$

Appendix

Mooney-Rivlin

It is a hyperelastic material model used for rubber-like materials as well as for biological tissues.

The *principal invariants* of $C := FF^T$ are defined as

$$I_1(C) := F \cdot F, \quad I_2(C) := F^* \cdot F^*$$

where $F^* := F^{-T} \det F$ is the *cofactor* of F .

References

- ▶ Nicola Pugno, Towards a Spiderman suit: large invisible cables and self-cleaning releasable superadhesive materials, *J. Phys.: Condens. Matter*, 19, 2007.
- ▶ Alessandro Granaldi, Paolo Decuzzi, The dynamic response of resistive microswitches: switching time and bouncing, *J. Micromech. Microeng.*, 16, 2006.
- ▶ Jiunn-Jong Wu, Adhesive contact between a nano-scale rigid sphere and an elastic half-space, *J. Phys. D: Appl. Phys.*, 39, 2006
- ▶ Z. J. Guo, N. E. McGruer, G. G. Adams, Modeling, simulation and measurement of the dynamic performance of an ohmic contact, electrostatically actuated RF MEMS switch, *J. Micromech. Microeng.*, 17, 2007
- ▶ Makoto Ashino, Alexander Schwarz, Hendrik Hölscher, Udo D. Schwarz, and Roland Wiesendanger, Interpretation of the atomic scale contrast obtained on graphite and single-walled carbon nanotubes in the dynamic mode of atomic force microscopy, *Nanotechnology*, 16, 2005

Supplementary references

- ▶ Gianfranco Capriz, Paolo Podio-Guidugli, Whence the boundary conditions in modern continuum physics?, *Atti Convegni Lincei n. 210*, 2004
- ▶ Antonio Di Carlo, Actual surfaces versus virtual cuts, *Atti Convegni Lincei n. 210*, 2004