Aneurismi sferici

Amabile Tatone

Roma 20 Gennaio 2007

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Hanno partecipato in particolare a questo progetto:

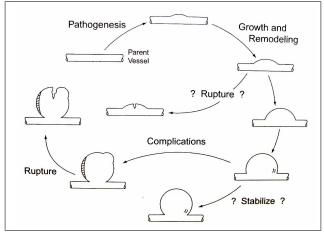
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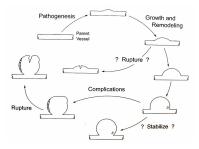
A possible natural history for saccular aneurysms

Intracranial saccular aneurysms are dilatations of the arterial wall.



[J.D.Humphrey, Cardiovascular Solid Mechanics, 2001]

A possible natural history for saccular aneurysms



- An initial insult may cause a local weakening of the wall and thus a mild dilatation.
- This raises the local stress field above normal values, thus setting into motion a growth and remodeling process that attempts to reduce the stress toward values that are homeostatic for the parent vessel.
- If degradation and deposition of collagen are well balanced, this could produce a larger, but stable lesion.
- If degradation exceedes deposition at any time, this could yield a rupture.

A recent survey in Japan

EBM of Neurosurgical Disease in Japan

Case No.	Sex	Age		Location of aneurysm	Motive for examination	Past history	Size of aneurysm (mm)				
							6 months	12 months	24 months	At rupture	Treatment
1	F	53	single	rt ICA	ex. for intracranial disease	hypertension	2.9	2.9	2.9	5.7	embolization (27 months)
2	F	71	multiple	AcomA	ex. for intracranial disease	hypertension, pituitary adenoma	4.9	4.9	4.9	5.9	clipping (7 months)
3	F	77	multiple	lt MCA	ex. for anxiety	hypertension, heart disease	4.5	4.5		4.5	dead (7 months)
4	F	42	multiple	lt MCA	ex. for anxiety	hypertension	4	4		7.0	clipping (18 months)

Table 2 Summary of patients with ruptured aneurysms

AcomA: anterior communicating artery, ex.: examination, ICA: internal carotid artery, MCA: middle cerebral artery.

	Clinical presentation								
	Time scale of d	levelopmer	nt	Multiple (148 aneurysms)	Single (232 aneurysms)	With SAH (30 aneurysms)	Total (380 aneurysms)		
				Mean follow up (months)					
	Days-Months Year		Decades	13.3	11.8	14.2	13.8		
Type 1		11 11	-	3 (2.0%)	1 (0.4%)		4 (1.0%)		
Type 3		₽		9 (6.1%)	9 (3.9%)	4 (13.3%)	18 (4.7%)		
Type 4			▲	136 (91.9%)	222 (95.7%)	26 (86.7%)) 358 (94.2%)		

Fig. 1 Process of growth and rupture of aneurysms. Type 1: aneurysm ruptures within a time span as short as several days to several months after formation. Type 2: aneurysm builds up slowly for a few years after formation and ruptures in this process, Type 3: aneurysm keeps growing slowly for many years without rupturing. Type 4: aneurysm grows up to a certain size, probably under 5

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Stroke

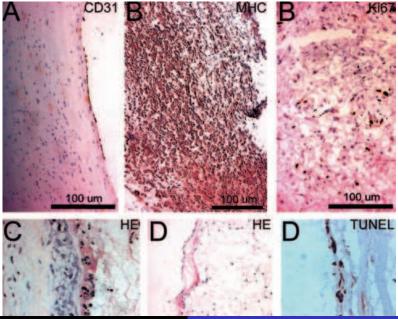
JOURNAL OF THE AMERICAN HEART ASSOCIATION

American Stroke Association

A Division of American Heart Association



Remodeling of Saccular Cerebral Artery Aneurysm Wall Is Associated With Rupture: Histological Analysis of 24 Unruptured and 42 Ruptured Cases Juhana Frösen, Anna Piippo, Anders Paetau, Marko Kangasniemi, Mika Niemelä, Juha Hernesniemi and Juha Jääskeläinen Stroke 2004;35;2287-2293; originally published online Aug 19, 2004;



Amabile Tatone

Aneurismi sferici

Before rupture, the wall of saccular cerebral artery aneurysms undergoes morphological changes associated with remodeling of the aneurysm wall. Some of these changes, like SMC [smooth muscle cell] proliferation and macrophage infiltration, likely reflect ongoing repair attempts that could be enhanced with pharmacological therapy. The morphological changes that result from the MH [myointimal hyperplasia] and matrix destruction are collectively referred to as remodeling of the vascular wall. Although MH is an adaptation mechanism of arteries to hemodynamic stress, in SAH [subarachnoid hemorrhage] patients, for undefined reasons, vascular wall remodeling [is] insufficient to prevent SCAA [saccular cerebral artery aneurysm] rupture.

- A. Di Carlo, S. Quiligotti, Growth and balance, *Mechanics Research Communications*, 29, 449–456, 2002.
- A. Di Carlo, Surface and bulk growth unified, *Mechanics of Material Forces* (P. Steinmann & G.A. Maugin, eds.), 53–64, Springer, New York, NY, 2005.
- M. Tringelová, P. Nardinocchi, L. Teresi and A. DiCarlo, The cardiovascular system as a smart system, in *Topics on Mathematics for Smart Systems*, eds. V. Valente and B. Miara (World Scientific, Singapore, 2007).

gross placement

$$p: \mathscr{B} \to \mathscr{E}$$

body gradient

$$\nabla p|_b : \mathrm{T}_b \mathscr{B} \to \mathrm{V} \mathscr{E}$$

element-wise configuration (prototype)

$$\mathbb{P}|_{b}: \mathrm{T}\mathscr{B}|_{b} \to \mathrm{V}\mathscr{E},$$

refined motion

$$(p(\tau), \mathbb{P}(\tau)) : \mathscr{B} \to \mathscr{E} \times (\mathcal{V}\mathscr{E} \otimes \mathcal{V}\mathscr{E})$$

realized velocity

$$(\dot{p}(\tau), \dot{\mathbb{P}}(\tau)\mathbb{P}(\tau)^{-1}): \mathscr{B} \to \mathcal{V}\mathscr{E} \times (\mathcal{V}\mathscr{E} \otimes \mathcal{V}\mathscr{E}).$$

test velocities

$$(\mathbf{v}, \mathbb{V}) : \mathscr{B} \to \mathcal{V}\mathscr{E} \times (\mathcal{V}\mathscr{E} \otimes \mathcal{V}\mathscr{E})$$

(gross velocity, growth velocity)

total working

$$\int_{\mathscr{B}} \Big(\, A^{i} \cdot \mathbb{V} - S \cdot Dv \, \Big) \; + \int_{\mathscr{B}} \big(\, b \cdot v + A^{\mathfrak{o}} \cdot \mathbb{V} \, \big) \; + \int_{\partial \mathscr{B}} t_{\partial \mathscr{B}} \cdot v$$

(integrals taken with respect to *prototypal volume* and *prototypal area*)

prototypal gradient

$$\operatorname{Dv} := (\nabla v) \mathbb{P}^{-1}$$

$$\int_{\mathscr{B}} \left(\operatorname{A}^{i} \cdot \mathbb{V} - \operatorname{S} \cdot \operatorname{Dv} \right) + \int_{\mathscr{B}} \left(\operatorname{b} \cdot \operatorname{v} + \operatorname{A}^{\mathfrak{o}} \cdot \mathbb{V} \right) + \int_{\partial \mathscr{B}} \operatorname{t}_{\partial \mathscr{B}} \cdot \operatorname{v} = \mathbf{0}$$

balance of brute forces

$$\begin{split} \mathrm{Div}\,\mathrm{S} + \mathrm{b} &= \mathsf{0} \ \ \mathrm{on} \ \mathscr{B} \\ \mathrm{S}\,\mathrm{n}_{\partial\mathscr{B}} &= \mathrm{t}_{\partial\mathscr{B}} \ \ \mathrm{on} \ \partial\mathscr{B} \end{split}$$

balance of accretive couples

$$A^{i} + A^{o} = 0$$
 on \mathscr{B}

free energy

$$\Psi(\mathscr{P}) \!=\! \int_{\mathscr{P}} \psi$$

(ψ free energy per unit prototypal volume)

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The constitutive theory of inner forces rests on two main pillars, altogether independent of balance:

- the principle of material indifference to change in observer
- the dissipation principle

Both of them deliver strict selection rules on admissible constitutive recipes for the inner force. None of them applies to the outer force, which has to be regarded as an adjustable control on the motion.

$$\int_{\mathscr{B}} \left(\,\mathrm{A}^i \cdot \mathbb{V} - \mathrm{S} \cdot \mathrm{D} \mathrm{v}\,\right) \; + \int_{\mathscr{B}} \left(\,\mathrm{b} \cdot \mathrm{v} + \mathrm{A}^{\mathfrak{o}} \cdot \mathbb{V}\,\right) \; + \int_{\partial \mathscr{B}} \mathrm{t}_{\partial \mathscr{B}} \cdot \mathrm{v} = \mathbf{0}$$

In our theory of the biomechanics of growth, the outer accretive couple A^o plays a primary role, representing the mechanical effects of the biochemical control system, smartly distributed within the body itself: ignoring the chemical degrees of freedom does not make negligible their feedback on mechanics.

Change in observer

$$\begin{split} \widetilde{p}(b,\tau) &= \widetilde{x}_{o}(\tau) + \widetilde{Q}(\tau) \left(p\left(b,\tau\right) - x_{o}(\tau) \right) \\ \widetilde{\mathbb{P}}(b,\tau) &= \mathbb{P}(b,\tau) \\ \widetilde{v}(b) &= \widetilde{Q}(\tau) v(b) + \widetilde{w}(\tau) + \widetilde{W}(\tau) \left(\widetilde{p}\left(b,\tau\right) - \widetilde{x}_{o}(\tau) \right) \\ \widetilde{\mathbb{V}}(b) &= \mathbb{V}(b) \end{split}$$

The working expended over each body-part on each test velocity (v, V) by the inner force constitutively related to each refined motion (p, P) should be invariant under all change in observer.

The free energy Ψ should be constitutively prescribed in such a way as to be invariant under all change in observer.

brute Cauchy stress

$$\mathrm{T} \mathrel{\mathop:}= (\mathsf{det}\,\mathrm{F})^{-1}\,\mathrm{S}\,\mathrm{F}^ op$$
 $\mathrm{T}^ op = \mathrm{T}$

warp

$$\mathrm{F} := \mathrm{D} p = (\nabla p) \mathbb{P}^{-1}$$

(measures how the body gradient of the gross placement differs from the prototypal stance)

Material indifference to change in observer

If we further assume that the response of the body element at b filters off from (p, \mathbb{P}) all information other than

 $p|_b, \nabla p|_b, \mathbb{P}|_b$

we obtain the following reduction theorem: there are constitutive mappings \widehat{S}_b , \widehat{A}_b^i and $\widehat{\psi}_b$ such that

$$S(b,\tau) = R(b,\tau) \widehat{S}_{b}(\ell_{b},\tau)$$

$$A^{i}(b,\tau) = \widehat{A}^{i}_{b}(\ell_{b},\tau)$$

$$\psi(b,\tau) = \widehat{\psi}_{b}(\ell_{b},\tau)$$

$$\ell_{b} := (U|_{b}, \mathbb{P}|_{b})$$

$$F = R U$$

$$\left(\mathbf{S}\cdot(\mathbf{D}\,\dot{\boldsymbol{\rho}})\,-\,\mathbf{A}^{\mathrm{i}}\cdot(\,\dot{\mathbb{P}}\mathbb{P}^{-1})\right)-\left(\dot{\psi}+\psi\,\mathbf{I}\cdot(\,\dot{\mathbb{P}}\mathbb{P}^{-1})\right)\geq\mathbf{0}$$

power expended : $-\{$ working expended by the inner force constitutively related to that motion on the velocity realized along the motion $\}$

power dissipated : {the power expended along a refined motion} {the time derivative of the free energy along that motion}

dissipation principle : the power dissipated should be non-negative, for all body-parts, at all times

$$\begin{split} \dot{\Psi}(\mathscr{P}) &= \left(\int_{\mathscr{P}} \psi \, \omega \right)^{\cdot} = \int_{\mathscr{P}} (\psi \, \omega)^{\cdot} \\ &= \int_{\mathscr{P}} (\dot{\psi} \, \omega \, + \, \psi \, \dot{\omega}) \\ &= \int_{\mathscr{P}} \left(\dot{\psi} + \psi \, \mathrm{I} \cdot (\, \dot{\mathbb{P}} \mathbb{P}^{-1}) \right) \omega \end{split}$$

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(the prototypal-volume form ω evolves in time with $\dot{\mathbb{P}}\mathbb{P}^{-1}$)

Constitutive assumptions: free energy and inner force

$$\psi(b,\tau) = \widehat{\psi}_{b}(\mathbf{U}|_{b}, \mathbb{P}|_{b}, \tau) = \varphi_{b}(\mathbf{F}(b,\tau))$$

The dissipation principle is fulfilled if and only if for each b the mappings \widehat{S} and \widehat{A}^i satisfy:

$$\widehat{\mathbf{S}} = \partial \varphi + \overset{+}{\mathbf{S}}, \qquad \widehat{\mathbf{A}}^{\mathfrak{i}} = \mathbb{E} + \overset{+}{\mathbf{A}}$$

Eshelby couple

$$\mathbb{E} \mathrel{\mathop:}= \mathbf{F}^{\!\!\top} \, \widehat{\mathbf{S}} - \varphi \, \mathbf{I}$$

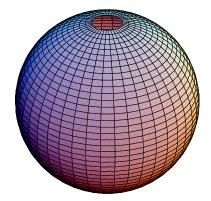
extra-energetic responses

$$\overset{\scriptscriptstyle +}{\mathrm{S}}, \overset{\scriptscriptstyle +}{\mathrm{A}}$$

reduced dissipation inequality

$$\overset{\mathrm{+}}{\mathrm{A}}(\ell, au)\cdot(\,\dot{\mathbb{P}}(au)\mathbb{P}(au)^{-1})-\overset{\mathrm{+}}{\mathrm{S}}(\ell, au)\cdot\dot{\mathrm{F}}(au)\,\leq\,0$$

Saccular aneurysms: Geometry & kinematics



Saccular aneurysms: Geometry & kinematics

paragon shape \mathcal{D} of \mathscr{B} centered at

$$x_{o} \in \mathscr{E}$$

spherical coordinates

$$(\widehat{\xi}(x),\widehat{\vartheta}(x),\widehat{\varphi}(x))$$

radius of x

$$\widehat{\xi}(x) = \|x - x_{o}\|$$

gross placement

$$p: \mathcal{D} \to \mathscr{E} x \mapsto x_{o} + \rho(\widehat{\xi}(x)) \operatorname{e}_{\mathsf{r}}(\widehat{\vartheta}(x), \widehat{\varphi}(x))$$

actual radius

$$\rho:[\xi_-,\xi_+]\to\mathbb{R}$$

spherically symmetric vector fields

$$v: \mathcal{D} \to V \mathscr{E}$$

radial component of v

 $\mathbf{v}: [\xi_{-}, \xi_{+}] \to \mathbb{R}$ $\mathbf{v}(\mathbf{x}) = \mathbf{v}(\xi) \mathbf{e}_{\mathbf{r}}(\vartheta, \varphi).$ spherically symmetric tensor fields

$$\begin{split} \mathscr{L}_{\varepsilon} &: \mathcal{D} \to \mathrm{V} \mathscr{E} \otimes \mathrm{V} \mathscr{E} \\ \mathscr{L}_{\varepsilon}(x) &= \mathrm{L}_{\mathsf{r}}(\xi) \operatorname{P}_{\mathsf{r}}(\vartheta, \varphi) + \mathrm{L}_{\mathsf{h}}(\xi) \operatorname{P}_{\mathsf{h}}(\vartheta, \varphi) \end{split}$$

orthogonal projectors

$$P_{\mathsf{r}}(x) := e_{\mathsf{r}}(\vartheta, \varphi) \otimes e_{\mathsf{r}}(\vartheta, \varphi)$$
$$P_{\mathsf{h}}(x) := I - P_{\mathsf{r}}(x)$$

∃ >

gradient of the gross placement

$$abla \mathbf{p}|_{\mathsf{x}} =
ho'(\xi) \operatorname{Pr}(\vartheta, \varphi) + rac{
ho(\xi)}{\xi} \operatorname{Ph}(\vartheta, \varphi),$$

prototype

$$\mathbb{P}(\mathbf{x},\tau) = \alpha_{\mathsf{r}}(\xi,\tau) \operatorname{P}_{\mathsf{r}}(\vartheta,\varphi) + \alpha_{\mathsf{h}}(\xi,\tau) \operatorname{P}_{\mathsf{h}}(\vartheta,\varphi) \,.$$

warp (Kröner-Lee decomposition)

$$\mathbf{F} := (\nabla \mathsf{p}) \, \mathbb{P}^{-1} = \, \lambda_{\mathsf{r}} \, \mathbf{P}_{\mathsf{r}} \, + \lambda_{\mathsf{h}} \mathbf{P}_{\mathsf{h}}$$

$$\lambda_{\mathsf{r}}(\xi,\tau) = \frac{\rho'(\xi,\tau)}{\alpha_{\mathsf{r}}(\xi,\tau)}, \qquad \lambda_{\mathsf{h}}(\xi,\tau) = \frac{\rho(\xi,\tau)}{\xi\alpha_{\mathsf{h}}(\xi,\tau)}$$

.

gross velocity and growth velocity

$$\dot{\mathsf{p}}(x,\tau) = \dot{\rho}(\xi,\tau) \operatorname{e}_{\mathsf{r}}(\vartheta,\varphi)$$
$$\dot{\mathbb{P}} \mathbb{P}^{-1}(x,\tau) = \frac{\dot{\alpha}_{\mathsf{r}}(\xi,\tau)}{\alpha_{\mathsf{r}}(\xi,\tau)} \operatorname{P}_{\mathsf{r}}(\vartheta,\varphi) + \frac{\dot{\alpha}_{\mathsf{h}}(\xi,\tau)}{\alpha_{\mathsf{h}}(\xi,\tau)} \operatorname{P}_{\mathsf{h}}(\vartheta,\varphi)$$

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working

(integrals taken with respect to the *paragon volume* and *paragon area*)

$$\begin{split} \int_{\xi_{-}}^{\xi_{+}} \left(\mathbf{A}_{\mathbf{r}} \mathbf{V}_{\mathbf{r}} + 2 \,\mathbf{A}_{\mathbf{h}} \mathbf{V}_{\mathbf{h}} - \mathbf{S}_{\mathbf{r}} \, \mathbf{v}' - 2 \,\mathbf{S}_{\mathbf{h}} \, \mathbf{v}/\xi \right) \mathbf{4} \, \pi \, \xi^{2} d\xi \, + \, \left(\mathbf{4} \, \pi \, \xi^{2} \, t \, \mathbf{v} \right) \Big|_{\xi_{\mp}} \\ & \mathbb{A} := \mathbb{A}^{\mathbf{i}} + \mathbb{A}^{\mathfrak{o}} = \mathbf{A}_{\mathbf{r}} \, \mathbf{P}_{\mathbf{r}} + \mathbf{A}_{\mathbf{h}} \, \mathbf{P}_{\mathbf{h}} \, . \end{split}$$

$$\begin{array}{lll} 2(S_{\mathsf{r}}(\xi) - S_{\mathsf{h}}(\xi)) + \xi S_{\mathsf{r}}'(\xi) &= 0 \\ & & & \\ A_{\mathsf{r}}(\xi) = A_{\mathsf{h}}(\xi) &= 0 \end{array} \right\} & (\xi_{-} < \xi < \xi_{+}) \\ & & & \\ & & \mp S_{\mathsf{r}}(\xi_{\mp}) = t_{\mp} \end{array}$$

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$$\Psi(\mathscr{P}) = \int_{\mathscr{P}} J\psi$$

(ψ free energy per unit *prototypal* volume)

$$J := \det(\mathbb{P}) = \alpha_{\mathsf{r}} \, \alpha_{\mathsf{h}}^2 > 0$$

($J\psi$ free energy per unit *paragon* volume)

Dissipation principle

$$\mathrm{S} \cdot (
abla \dot{\mathsf{p}}) - \mathbb{A}^{\mathrm{i}} \cdot (\ \dot{\mathbb{P}} \ \mathbb{P}^{-1}) - (J \ \psi)^{\cdot} \geq 0$$

$$\psi(\mathbf{x},\tau) = \varphi\left(\lambda_{\mathsf{r}}(\xi,\tau),\lambda_{\mathsf{h}}(\xi,\tau);\xi\right)$$

$$\begin{split} \mathbf{S}_{\mathsf{r}} &= J\varphi_{,\mathsf{r}}/\alpha_{\mathsf{r}} + \overset{+}{\mathbf{S}}_{\mathsf{r}} & \mathbf{S}_{\mathsf{h}} &= J\varphi_{,\mathsf{h}}/\alpha_{\mathsf{h}} + \overset{+}{\mathbf{S}}_{\mathsf{h}} \\ \mathbf{A}_{\mathsf{r}}^{\mathsf{i}} &= J\left[\mathbf{S}_{\mathsf{r}}\,\alpha_{\mathsf{r}}\,\lambda_{\mathsf{r}}/J - \phi\right] + \overset{+}{\mathbf{A}}_{\mathsf{r}} & \mathbf{A}_{\mathsf{h}}^{\mathsf{i}} &= J\left[\mathbf{S}_{\mathsf{h}}\,\alpha_{\mathsf{h}}\,\lambda_{\mathsf{h}}/J - \phi\right] + \overset{+}{\mathbf{A}}_{\mathsf{h}} \end{split}$$

reduced dissipation inequality

$$\overset{\scriptscriptstyle +}{\mathrm{Sr}} \alpha_{\mathsf{r}} \, \dot{\lambda}_{\mathsf{r}} + 2 \overset{\scriptscriptstyle +}{\mathrm{Sh}} \alpha_{\mathsf{h}} \, \dot{\lambda}_{\mathsf{h}} - \overset{\scriptscriptstyle +}{\mathrm{A}_{\mathsf{r}}} \, \dot{\alpha}_{\mathsf{r}} / \alpha_{\mathsf{r}} - 2 \overset{\scriptscriptstyle +}{\mathrm{A}_{\mathsf{h}}} \, \dot{\alpha}_{\mathsf{h}} / \alpha_{\mathsf{h}} \ge 0$$

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$$\overset{+}{\mathrm{S}}_{\mathsf{r}}\,\alpha_{\mathsf{r}}\,\dot{\lambda}_{\mathsf{r}}+2\overset{+}{\mathrm{S}}_{\mathsf{h}}\,\alpha_{\mathsf{h}}\,\dot{\lambda}_{\mathsf{h}}-\overset{+}{\mathrm{A}}_{\mathsf{r}}\,\dot{\alpha}_{\mathsf{r}}/\alpha_{\mathsf{r}}-2\overset{+}{\mathrm{A}}_{\mathsf{h}}\,\dot{\alpha}_{\mathsf{h}}/\alpha_{\mathsf{h}}\geq 0$$

$$\begin{split} \dot{\vec{S}}_{r} &= \dot{\vec{S}}_{h} = 0 \\ \dot{\vec{A}}_{r} &= -JD_{r} \, \dot{\alpha}_{r} \,, \qquad \dot{\vec{A}}_{h} = -JD_{h} \, \dot{\alpha}_{h} / \alpha_{h} \\ D_{r} &> 0 \,, \ D_{h} > 0 \end{split}$$

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incompressibility

$$\det \mathrm{F} = \lambda_\mathsf{r}\,\lambda_\mathsf{h}^2 = 1 \quad \Longleftrightarrow \quad \lambda_\mathsf{r} = 1/\lambda_\mathsf{h}^2\,.$$

reactive inner force

$$\overset{\bowtie}{\mathbf{S}} = J^{\bowtie}_{\pi} \left(\frac{1}{\alpha_{\mathsf{r}} \, \lambda_{\mathsf{r}}} \, \mathbf{P}_{\mathsf{r}} + \frac{1}{\alpha_{\mathsf{h}} \, \lambda_{\mathsf{h}}} \, \mathbf{P}_{\mathsf{h}} \right), \qquad \overset{\bowtie}{\mathbb{A}} = J^{\bowtie}_{\pi} \, I$$

free-energy restriction

$$\widetilde{\varphi}: \lambda \mapsto \varphi(1/\lambda^2, \lambda)$$

active and reactive components

$$\begin{split} \mathbf{S}_{\mathsf{r}} &= \frac{J}{\alpha_{\mathsf{r}}\,\lambda_{\mathsf{r}}} \left(\overset{\bowtie}{\pi} - \left(\lambda_{\mathsf{h}}/3 \right) \widetilde{\varphi}^{\,\prime} \right) \qquad \qquad \mathbf{S}_{\mathsf{h}} &= \frac{J}{\alpha_{\mathsf{h}}\,\lambda_{\mathsf{h}}} \left(\overset{\bowtie}{\pi} + \left(\lambda_{\mathsf{h}}/6 \right) \widetilde{\varphi}^{\,\prime} \right) \\ \mathbf{A}_{\mathsf{r}}^{\mathsf{i}} &= J \left(\mathbf{T}_{\mathsf{r}} - \widetilde{\varphi} - D_{\mathsf{r}}\,\dot{\alpha}_{\mathsf{r}}/\alpha_{\mathsf{r}} \right) \qquad \qquad \mathbf{A}_{\mathsf{h}}^{\mathsf{i}} &= J \left(\mathbf{T}_{\mathsf{h}} - \widetilde{\varphi} - D_{\mathsf{h}}\,\dot{\alpha}_{\mathsf{h}}/\alpha_{\mathsf{h}} \right) \end{split}$$

Cauchy stress

$$T = (J \operatorname{\mathsf{det}}(F))^{-1} \operatorname{S} \mathbb{P}^\top F^\top$$

radial and hoop components

$$T_{\rm r} = J^{-1} S_{\rm r} \, \alpha_{\rm r} \, \lambda_{\rm r} \tag{1}$$

$$T_{h} = J^{-1}S_{h} \alpha_{h} \lambda_{h}$$
(2)

$$\widetilde{arphi}(\lambda) = (c/\delta) \exp((\Gamma/2) (\lambda^2 - 1)^2),$$

 $c = 0.88 N/m, \quad \Gamma = 12.99$

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Constitutive recipes for the outer accretive couple

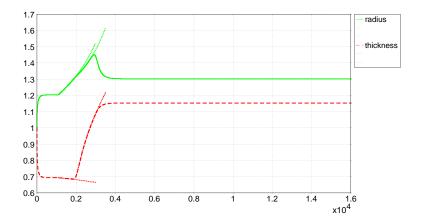
$$\begin{split} \mathrm{A}_{\mathsf{r}}^{\mathfrak{o}} &= J\left(\mathit{G}_{\mathsf{r}}\left(\mathrm{T}_{\mathsf{h}}-\mathrm{T}^{\odot}\right)-\mathrm{T}_{\mathsf{r}}+\widetilde{\varphi}\right),\\ \mathrm{A}_{\mathsf{h}}^{\mathfrak{o}} &= J\left(\mathit{G}_{\mathsf{h}}\left(\mathrm{T}^{\odot}-\mathrm{T}_{\mathsf{h}}\right)-\mathrm{T}_{\mathsf{h}}+\widetilde{\varphi}\right), \end{split}$$

control gains

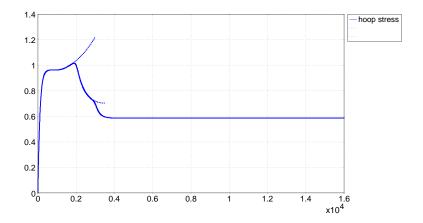
 $G_{\rm r}, G_{\rm h}$

Evolution laws

$$\begin{split} \dot{\alpha}_{\rm r}/\alpha_{\rm r} \;&=\; ({\it G}_{\rm r}/{\it D}_{\rm r}) \left({\rm T}_{\rm h}-{\rm T}^\odot\right), \\ \dot{\alpha}_{\rm h}/\alpha_{\rm h} \;&=\; ({\it G}_{\rm h}/{\it D}_{\rm h}) \left({\rm T}^\odot-{\rm T}_{\rm h}\right). \end{split}$$

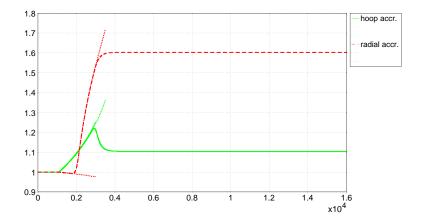


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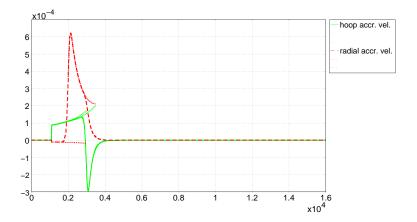
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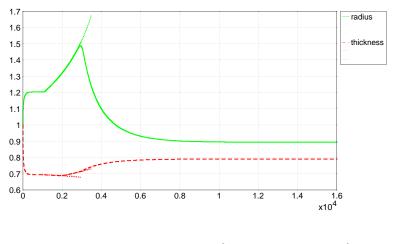


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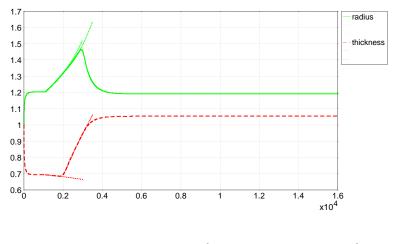
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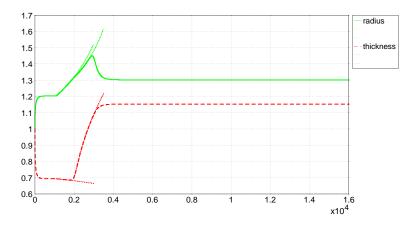


 $D_{\rm r}/D_{\rm h} = 1$ $G_{\rm r}/D_{\rm r} = 10^2$ $G_{\rm h}/D_{\rm h} = 10^3$



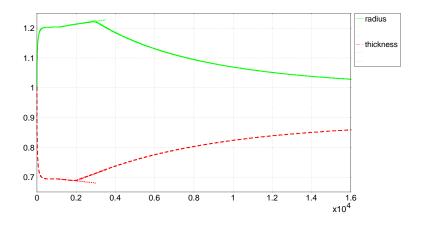
 $D_r/D_h = 1$ $G_r/D_r = 10^3$ $G_h/D_h = 2.5 \times 10^3$

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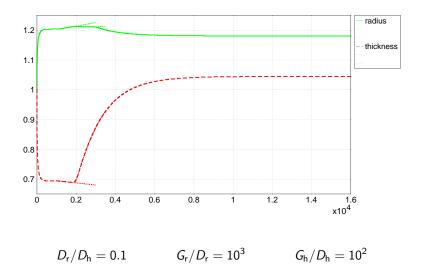


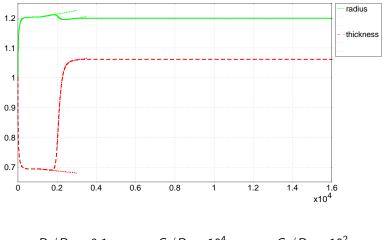
 $D_r/D_h = 1$ $G_r/D_r = 2 \times 10^3$ $G_h/D_h = 5 \times 10^3$

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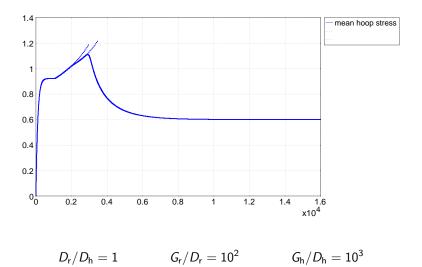


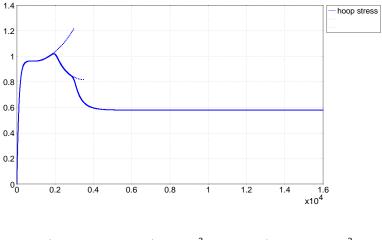
 $D_{\rm r}/D_{\rm h} = 0.1$ $G_{\rm r}/D_{\rm r} = 10^2$ $G_{\rm h}/D_{\rm h} = 10^2$



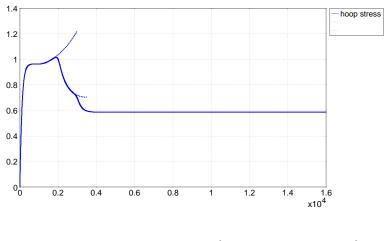


 $D_r/D_h = 0.1$ $G_r/D_r = 10^4$ $G_h/D_h = 10^2$

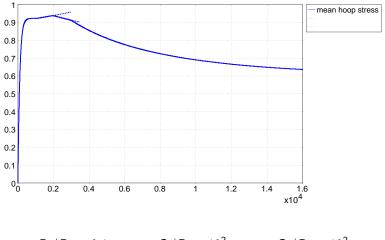




 $D_{\rm r}/D_{\rm h} = 1$ $G_{\rm r}/D_{\rm r} = 10^3$ $G_{\rm h}/D_{\rm h} = 2.5 \times 10^3$

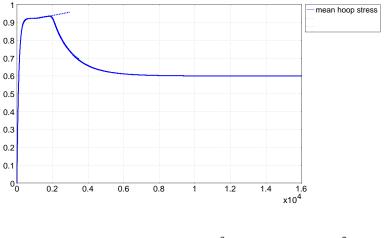


 $D_{\rm r}/D_{\rm h} = 1 \qquad \quad G_{\rm r}/D_{\rm r} = 2 \times 10^3 \qquad \quad G_{\rm h}/D_{\rm h} = 5 \times 10^3$



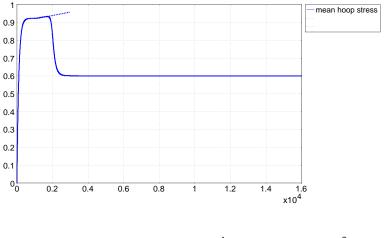
 $D_r/D_h = 0.1$ $G_r/D_r = 10^2$ $G_h/D_h = 10^2$

<ロ> <四> <四> <四> <三</td>



 $D_{\rm r}/D_{\rm h} = 0.1$ $G_{\rm r}/D_{\rm r} = 10^3$ $G_{\rm h}/D_{\rm h} = 10^2$

<ロ> <四> <四> <四> <三</td>



 $D_r/D_h = 0.1$ $G_r/D_r = 10^4$ $G_h/D_h = 10^2$

<ロ> <四> <四> <四> <三</td>