Homeostatic states and relaxed configurations of a spherical thick shell

Amabile Tatone

Dipartimento di Ingegneria delle Strutture, delle Acque e del Terreno Università dell'Aquila - Italy

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A joint work with:

A. Di Carlo^{*a*}, V. Varano^{*a*}.

^a Università degli Studi "Roma Tre", Roma

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Kröner-Lee scheme



$$\mathbf{F} := \nabla \mathbf{p} \, \mathbf{G}^{-1}$$

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Transplacement

$$\mathbf{p}(\mathbf{x},\tau) \,=\, \mathbf{x}_{\mathsf{o}} + \rho\left(\xi,\tau\right) \mathbf{e}_{\mathsf{r}}(\vartheta,\varphi)\,,$$

Transplacement gradient and remodeling

$$\nabla \mathbf{p}|_{\mathbf{x}} = \rho'(\xi) \, \mathbf{P}_r + \xi^{-1} \rho(\xi) \, \mathbf{P}_h$$
$$\mathbf{G}(\mathbf{x}) = \gamma_r(\xi) \, \mathbf{P}_r + \gamma_h(\xi) \, \mathbf{P}_h$$

Effective stretch

$$\mathbf{F}(\mathbf{x}) := \nabla \mathbf{p}|_{\mathbf{x}} \mathbf{G}(\mathbf{x})^{-1} = \lambda_r(\xi) \, \mathbf{P}_r + \lambda_h(\xi) \, \mathbf{P}_h$$

Effective stretches

$$\lambda_h(\xi) := \frac{\rho(\xi)}{\xi \gamma_h(\xi)}, \quad \lambda_r(\xi) := \frac{\rho'(\xi)}{\gamma_r(\xi)}$$

Jacobian determinant

$$J_{\rho}(\xi) := (\rho(\xi)/\xi)^2 \, \rho'(\xi)$$

Relaxed Jacobian

$$J(\xi) := \gamma_h^2(\xi) \gamma_r(\xi)$$

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Balance

$$2\left(S_r(\xi) - S_h(\xi)\right) + \xi S'_r(\xi) = 0$$
$$S_r(\xi_{\mp}) = -\pi_{\mp}$$

 π_{\mp} outer *reference* pressure.

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Cauchy stress

$$T_h(\xi) = \frac{\rho(\xi)}{\xi} S_h(\xi) / J_p(\xi) , \quad T_r(\xi) = \rho'(\xi) S_r(\xi) / J_p(\xi)$$

Balance equations

$$\frac{\rho(\xi)}{\xi} \left(2 \left(\mathrm{T}_r(\xi) - \mathrm{T}_h(\xi) \right) \rho'(\xi) + \rho(\xi) \, \mathrm{T}'_r(\xi) \right) = 0$$
$$\mathrm{T}_r(\xi_{\mp}) = -\,\widehat{p}_{\mp}$$

 \widehat{p}_{\mp} outer *actual* pressure

$$\pi_{\mp} = \left(\rho\left(\xi_{\mp}\right)/\xi_{\mp}\right)^2 \widehat{p}_{\mp}$$

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Strain energy function

Elastically incompressible material

$$\lambda_h^2(\xi)\lambda_r(\xi) = 1 \tag{1}$$

Fung strain energy function

$$\phi(\lambda_h) = \frac{c}{H} \left(-1 + e^{\frac{1}{2}\Gamma\left(\lambda_h^2 - 1\right)^2} \right)$$
(2)

Because of the incompressibility constraint the response function turns out to be defined only for the deviatoric part of the stress:

$$T_h(\xi) - T_r(\xi) = \frac{\lambda_h(\xi)}{2} \phi'(\lambda_h(\xi))$$
(3)

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For a given shape and fixed values of the outer pressure, find the stress and the transformation stretch. It is assumed that the hoop stress field is uniform across the thickness.

Results: It is shown that there exists only one solution (although no formal proof is given). The main point is that the transplacement is the identity map. This also makes Piola stress and Cauchy stress indistinguishable. The reference shape is the same as the given shape:

$$\rho(\xi) = \xi \quad \Rightarrow \quad \lambda_h \gamma_h = 1, \ \lambda_r \gamma_r = 1$$

We assume that a homeostatic state is characterized by a uniform hoop stress S_h^{\diamond} . The balance equation gives the following solution for the radial stress

$$\mathrm{S}_r(\xi) = \mathrm{S}_h^\diamond + \frac{C}{\xi^2}$$

By enforcing the two boundary conditions

$$\mathbf{S}_r(\xi_{\mp}) = -\pi_{\mp}$$

we get both C and S_h^\diamond

Piola stress field

$$S_r(\xi) = \frac{\pi_+ \left(\xi^2 - \xi_-^2\right)\xi_+^2 + \pi_- \left(\xi_+^2 - \xi^2\right)\xi_-^2}{\xi^2 \left(\xi_-^2 - \xi_+^2\right)}$$

$$\mathbf{S}_{h}^{\diamond} = \frac{\pi_{+}\,\xi_{+}^{2} - \pi_{-}\,\xi_{-}^{2}}{\left(\xi_{-}^{2} - \xi_{+}^{2}\right)}$$

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Numerical values

$$\begin{split} \xi_m &= 2000/20 \ \mu\text{m} \\ L &= H/1.18^2 \ \mu\text{m} \\ c &= 0.8769 \times 10^{-6} \times 10^{-1} \ \text{N} \ \mu\text{m}^{-2} = 876.9 \times 10^{-1} \ \text{kPa} \\ \Gamma &= 12.99 \\ H &= 27.8 \ \mu\text{m} \\ \pi_+ &= 2 \times 10^{-9} \ \text{N} \ \mu\text{m}^{-2} = 2 \ \text{kPa} \\ \pi_- &= 12 \times 10^{-9} \ \text{N} \ \mu\text{m}^{-2} = 12 \ \text{kPa} \end{split}$$

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Initial homeostatic state





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Effective stretch



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Transformation stretch



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Strain energy density



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After setting the outer pressure to zero, let the body relax while keeping the transformation stretch unchanged.

Results: The configuration the body reaches is not stress free.

Unloading transplacement

Since both γ_h and γ_r are left unchanged the transplacement is isochoric. From

$$\lambda_h(\xi) = \frac{\rho(\xi)}{\xi \gamma_h(\xi)}, \quad \lambda_r(\xi) = \frac{\rho'(\xi)}{\gamma_r(\xi)}, \quad \lambda_h^2(\xi)\lambda_r(\xi) = 1$$

we get

$$\rho(\xi) = \left(\xi^3 + (\rho_m^3 - \xi_m^3)\right)^{1/3} \tag{4}$$

where $\rho_m := \rho(\xi_m)$ is an integration constant.

From material response and balance equation

$$T_h(\xi) - T_r(\xi) = \frac{\lambda_h(\xi)}{2} \phi'(\lambda_h(\xi))$$
$$2 \left(T_r(\xi) - T_h(\xi)\right) \rho'(\xi) + \rho(\xi) T'_r(\xi) = 0$$

we get the boundary value problem

$$-\lambda_h(\xi) \phi'(\lambda_h(\xi)) \rho'(\xi) + \rho(\xi) \operatorname{T}'_r(\xi) = 0$$
$$\operatorname{T}_r(\xi_{\mp}) = 0$$

which we can solve (possibly numerically) for both T_r and ρ_m , by using also $\lambda_h(\xi) := \rho(\xi)/\xi \gamma_h(\xi)$. Finally we compute T_h .

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Residual effective stretch



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Residual effective stretch



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Residual stress



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Residual stress



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Residual strain energy density



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Residual strain energy density



Relaxed configuration

Does a relaxed configuration exist?

$$\nabla \mathbf{p}|_{\mathbf{x}} = \rho'(\xi) \, \mathbf{P}_r + \xi^{-1} \rho(\xi) \, \mathbf{P}_h$$
$$\lambda_r(\xi) = \frac{\rho'(\xi)}{\gamma_r(\xi)}, \quad \lambda_h(\xi) = \frac{\rho(\xi)}{\xi \gamma_h(\xi)}$$
$$\lambda_r(\xi) = 1 \& \lambda_h(\xi) = 1 \implies \gamma_r(\xi) - (\gamma_h(\xi) + \xi \gamma'_h(\xi)) = 0$$



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Kröner-Lee scheme



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Unloaded two layer shape



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Unloaded two layer shape



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Residual effective stretch



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Residual effective stretch



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Residual stress



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Residual stress



Residual strain energy density



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Residual strain energy density



Reference shape



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Slice

Unloaded two layer shape



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Slice

Unloaded two layer shape





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Open slice split into 2 layers



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Open slice split into 2 layers



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Open slice split into 4 layers



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Open slice split into 6 layers



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