

## Cauchy stress tensor and Piola stress tensor

$T \cdot \nabla_N$  stress power density  
per unit current volume

$\nabla_N F = \nabla_{\bar{r}}$  velocity gradient in the reference shape

$$\begin{aligned}
 \underbrace{(T \cdot \nabla_N)}_{\text{CAUCHY STRESS}} V_{\bar{r}} &= (T \cdot \nabla_N) (\det F) V_{\bar{r}} \\
 &= (\det F) (T \cdot \nabla_{\bar{r}} F^{-1}) V_{\bar{r}} \\
 &= \underbrace{(\det F) (T F^{-T})}_S \cdot \nabla_{\bar{r}} V_{\bar{r}} = \underbrace{S \cdot \nabla_{\bar{r}}}_{\text{PIOLA STRESS}} V_{\bar{r}}
 \end{aligned}$$

$\text{cof } F := (\det F) F^{-T}$  cofactor of  $F$

$$\begin{aligned}
 \int_{\bar{\Omega}}^{\text{int}} (r) &= - \int_{\bar{\Omega}} T \cdot \nabla_N dV = - \int_{\bar{\Omega}} (T \cdot \nabla_N) \det F dV \\
 &= - \int_{\bar{\Omega}} S \cdot \nabla_{\bar{r}} dV
 \end{aligned}$$

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Force balance (power balance)

$$\mathcal{J}^{\text{ext}}(\mathcal{R}) + \mathcal{J}^{\text{int}}(\mathcal{R}) = 0 \quad \forall \mathcal{R}$$

Energy imbalance principle

$$\mathcal{J}^{\text{ext}}(\mathcal{R}) - \frac{d}{dt} \int_{\bar{\mathcal{R}}} \varphi(F) dV \geq 0$$

in any motion.

$\varphi(F)$  strain energy density per unit reference volume

Since

$$\mathcal{J}^{\text{ext}}(\mathcal{R}) = -\mathcal{J}^{\text{int}}(\mathcal{R}) = \int_{\mathcal{R}} T \cdot \dot{F} F^{-1} dV = \int_{\bar{\mathcal{R}}} S \cdot \dot{F} dV$$

we get

$$\int_{\bar{\mathcal{R}}} S \cdot \dot{F} dV - \frac{d}{dt} \int_{\bar{\mathcal{R}}} \varphi(F) dV \geq 0$$

We extend this inequality to any subset of  $\bar{\mathcal{R}}$  and get

$$S \cdot \dot{F} - \frac{d}{dt} \varphi(F) \geq 0$$

If  $S = \hat{S}(F)$ , and correspondingly  $T = \hat{T}(F)$ ,  
the material is an elastic material with  
response function  $\hat{S}(F)$ , and correspondingly  $\hat{T}(F)$ .

[2018-04-25]

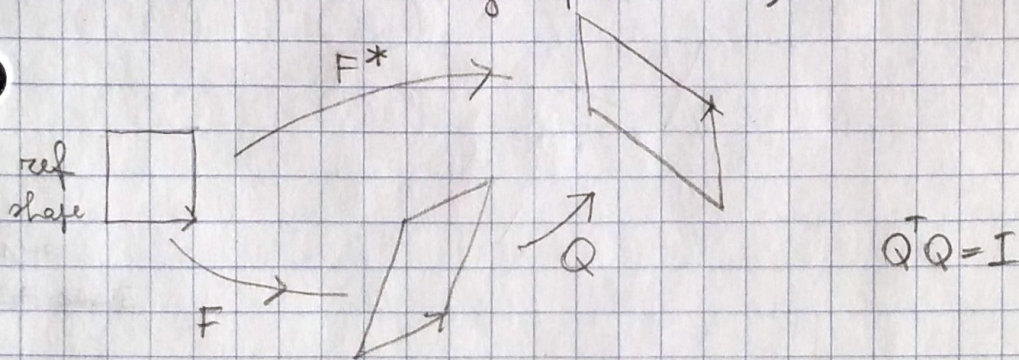
## Strain energy

If

$$\hat{S}(F) \cdot \dot{F} = \frac{d}{dt} \varphi(F)$$

the material is called a hyperelastic material.

We require  $\varphi$  to be frame indifferent (objective, invariant under a change of observer)



$$\varphi(F) = \varphi(F^*) \quad F^* = QF$$

$$\Rightarrow \varphi(F) = \varphi(QF) \quad \forall Q, \forall F$$

$$\varphi(F) = \varphi(QRU)$$

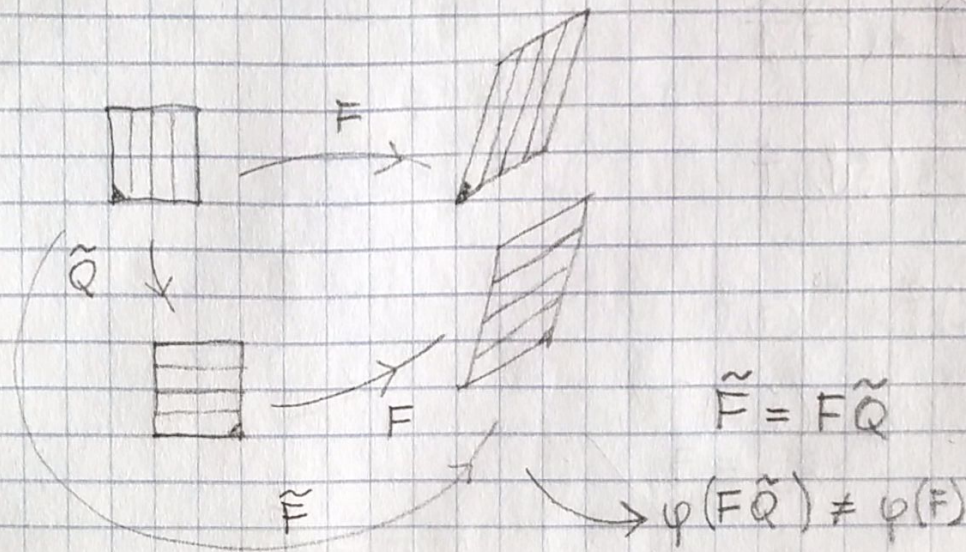
$$\varphi(F) = \varphi(RU) = \varphi(R^*U) \quad R^* = QR$$

$\Rightarrow \varphi$  is independent of the rotation tensor

$$\varphi(F) = \varphi(U)$$

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## Material symmetry



$\tilde{Q}$  symmetry

$\varphi(F\tilde{Q}) = \varphi(F)$  symmetry invariance

$\varphi(RU\tilde{Q}) = \varphi(U)$  obj

$\varphi(R\tilde{Q}\tilde{Q}^T U \tilde{Q})$

obj  $\varphi(\tilde{Q}^T U \tilde{Q}) = \varphi(U)$

Symmetry group  $G = \{\tilde{Q}_1, \tilde{Q}_2, \dots, I\} \subset \text{Orth}^+$

for a material; it characterizes both a material and its reference shapes with a zero stress.

If  $G = \text{Orth}^+$  (the rotation group)

the material is called isotropic material.