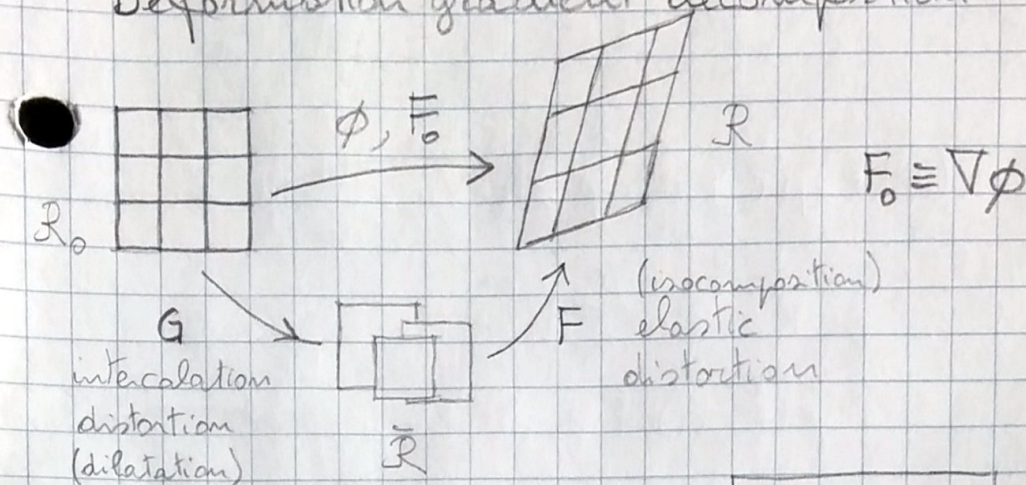


## Deformation gradient decomposition



$$\beta = \det \mathbb{G} \quad \mathbb{G} = \beta^{1/3} \mathbf{I} \quad \boxed{\mathbb{F}_0 = \mathbb{F} \mathbb{G}}$$

$$\det \mathbb{F}_0 = \beta (\det \mathbb{F})$$

 stoichiometric constant  $\downarrow$ 

$$\det \mathbb{G} = (\rho_0 V_0 + \rho_L V_L) / (\rho_0 V_0) = 1 + c \frac{V_L}{V_0} = 1 + c \alpha$$

$$\dot{\beta} = \alpha \dot{c} \quad \frac{d}{dt} (\det \mathbb{G}) = (\det \mathbb{G}) \operatorname{tr} (\dot{\mathbb{G}} \mathbb{G}^{-1}) = \dots = \alpha \dot{c}$$

$$\dot{\mathbb{G}} \mathbb{G}^{-1} = \frac{1}{3} \beta^{-2/3} \beta^{-1/3} \alpha \dot{c} \mathbf{I} = \frac{1}{3} \frac{\alpha}{\beta} \dot{c} \mathbf{I}$$

$$\dot{\mathbb{F}}_0 \mathbb{F}_0^{-1} = (\dot{\mathbb{F}} \mathbb{G} + \mathbb{F} \dot{\mathbb{G}}) \mathbb{G}^{-1} \mathbb{F}^{-1} = \dot{\mathbb{F}} \mathbb{F}^{-1} + \frac{1}{3} \frac{\alpha}{\beta} \dot{c} \mathbf{I}$$

$$\mathbb{S}_0 = \mathbf{T} \mathbb{F}_0^{-T} (\det \mathbb{F}_0) \quad \mathbb{S} = \mathbf{T} \mathbb{F}^{-T} (\det \mathbb{F})$$

$$\mathbb{S}_0 \cdot \dot{\mathbb{F}}_0 = (\det \mathbb{F}_0) \mathbf{T} \mathbb{F}_0^{-T} \cdot \dot{\mathbb{F}}_0 = \beta (\det \mathbb{F}) \mathbf{T} \cdot \dot{\mathbb{F}} \mathbb{F}^{-1}$$

$$= \beta (\det \mathbb{F}) \mathbf{T} \cdot \dot{\mathbb{F}} \mathbb{F}^{-1} + \frac{1}{3} \alpha (\det \mathbb{F}) \mathbf{T} \cdot \mathbf{I} \dot{c} = \beta \mathbb{S} \cdot \dot{\mathbb{F}} - \alpha (\det \mathbb{F}) p \dot{c}$$

$$\text{with } -p = \frac{1}{3} \mathbf{T} \cdot \mathbf{I} = \frac{1}{3} \operatorname{tr} \mathbf{T}$$

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Hence the energy imbalance takes the form

$$\beta S \cdot \dot{F} - \alpha (\det F) \rho \dot{c} + \mu_0 \dot{c} \rho_0 - h_0 \cdot \nabla \mu_0 - \frac{d}{dt} \psi \geq 0$$

If we choose for the free energy the expression

$$\hat{\psi}(F, c) = \beta \varphi(F)$$

while assuming

$$\frac{d}{dt} \varphi(F) = \hat{S}(F) \cdot \dot{F}$$

then

$$\frac{d}{dt} \psi = \dot{\beta} \varphi(F) + \beta \hat{S}(F) \cdot \dot{F} = \alpha \dot{c} \varphi(F) + \beta \hat{S}(F) \cdot \dot{F}$$

We finally get

$$\beta (S - \hat{S}(F)) \cdot \dot{F} - \alpha \dot{c} \varphi(F) - (\det F) \rho \alpha \dot{c} + \mu_0 \dot{c} \rho_0 - h_0 \cdot \nabla \mu_0 \geq 0$$

which, rearranging terms, becomes

$$\beta \underbrace{(S - \hat{S}(F))}_{S^+} + \underbrace{\left( \frac{\rho_0}{\alpha} \mu_0 - (\det F) \rho - \varphi(F) \right)}_{\mu^+} \alpha \dot{c} - h_0 \cdot \nabla \mu_0 \geq 0$$

$$\Rightarrow \begin{cases} S = \hat{S}(F) + S^+ & S^+ \cdot \dot{F} \geq 0 \\ \mu_0 = \frac{\alpha}{\rho_0} \left( (\det F) \rho + \varphi(F) \right) + \mu^+ & \mu^+ \cdot \dot{c} \geq 0 \\ -h_0 \cdot \nabla \mu_0 \geq 0 \end{cases}$$



Let us extend the free energy expression introducing an energy term depending just on the concentration

$$\hat{\psi}(\mathbf{F}_0, c) = \varphi_{ch}(c) + \beta \varphi_e(\mathbf{F})$$

and setting  $\frac{d}{dt} \varphi_{ch}(c) = \rho_0 \hat{\mu}_{ch}(c) \dot{c}$

We get

$$\begin{aligned} \frac{d}{dt} \hat{\psi}(\mathbf{F}_0, c) &= \rho_0 \hat{\mu}_{ch}(c) \dot{c} + \dot{\beta} \varphi_e(\mathbf{F}) + \beta \hat{\mathbf{S}}(\mathbf{F}) \cdot \dot{\mathbf{F}} \\ &= (\rho_0 \hat{\mu}_{ch}(c) + \alpha \varphi_e(\mathbf{F})) \dot{c} + \beta \hat{\mathbf{S}}(\mathbf{F}) \cdot \dot{\mathbf{F}} \end{aligned}$$

From the energy imbalance expression we arrive at

$$\mu_0 = \hat{\mu}_{ch}(c) + \frac{\alpha}{\rho_0} \left( (\det \mathbf{F}) p + \varphi_e(\mathbf{F}) \right) + \mu^*$$

with

$$p = -\frac{1}{3} \text{tr} \mathbf{T}$$



Summarizing

$$S = \hat{S}_e(F) + S^+$$

$$S^+ = (\det F) T^+ F^{-T}$$

$$T = \hat{T}_e(F) + T^+$$

$$T^+ = 2\mu \operatorname{sym}(\dot{F}F^{-1})$$

For an almost incompressible neo-Hookean material

$$\hat{T}_e(F) = \underbrace{2k_I J^{-\frac{5}{3}} \left( B - \frac{1}{3} I_1 I \right)}_{\text{deviatoric}} + \underbrace{2k_V (J-1)}_{\text{spherical}}$$

$$\hat{S}_e(F) = 2k_I J^{-\frac{2}{3}} \left( F - \frac{1}{3} I_1 F^{-T} \right) + 2k_V (J-1) J F^{-T}$$

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